

## Magnetic flux $\Phi$

$$\Phi = \int_S \vec{B} \cdot d\vec{s} = \oint_C \vec{A} \cdot d\vec{l} \quad (\text{wb})$$

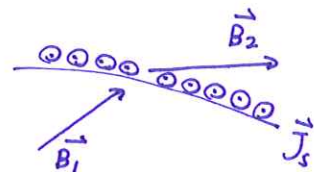
## Magnetic Permeability

$$\vec{B} = \mu_0 (\vec{H} + \vec{M}) = \mu_0 (\vec{H} + \chi_m \vec{H}) = \mu_0 (1 + \chi_m) \vec{H}$$

$$\vec{B} = \mu_0 \mu_r \vec{H}$$

## Magnetic Boundary Condition

$$\begin{cases} B_{1n} = B_{2n} \\ H_{2t} - H_{1t} = J_s \quad (\frac{A}{m}) \end{cases}$$



Self-inductance:  $L \equiv \frac{\Lambda}{I} \quad (\text{H})$      $\Lambda$ : magnetic flux linkage =  $N \int \vec{B} \cdot d\vec{s}$

For a solenoid:  $L = \mu \frac{N^2}{l} S$



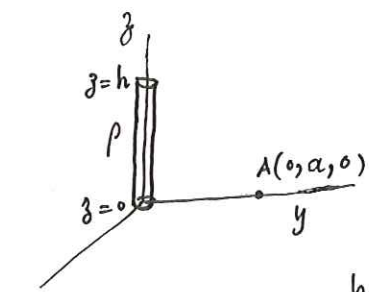
Mutual inductance:

$$L_{12} = \frac{\Lambda_{12}}{I_1} = \frac{N_2}{I_1} \int_{S_2} \vec{B}_1 \cdot d\vec{s} \quad (\text{H})$$

Magnetic Energy:

$$W_m = \frac{1}{2} \int_V \mu H^2 dv$$

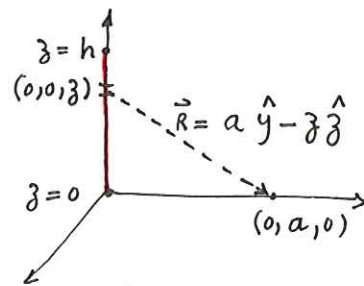
Example 1: For a line of charge as depicted in the picture calculate the electric field at point A.



Solution:  $\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\rho dl}{R^3} \vec{R}$

$$\vec{R} = a\hat{y} - z\hat{z}$$

$$R = (a^2 + z^2)^{1/2}$$



$$\begin{aligned} \vec{E} &= \frac{\rho}{4\pi\epsilon_0} \int_0^h \frac{dz}{(a^2 + z^2)^{3/2}} (a\hat{y} - z\hat{z}) = \frac{\rho}{4\pi\epsilon_0} \left[ \hat{y} \int_0^h \frac{adz}{(a^2 + z^2)^{3/2}} - \hat{z} \int_0^h \frac{zdz}{(a^2 + z^2)^{3/2}} \right] \\ &= \frac{\rho}{4\pi\epsilon_0} \left[ \hat{y} \int_0^h \frac{a^2 + z^2 - z^2}{a(a^2 + z^2)^{3/2}} dz - \hat{z} \frac{1}{2} \int_0^h \frac{2zdz}{(a^2 + z^2)^{3/2}} \right] \end{aligned}$$

$$= \frac{\rho}{4\pi\epsilon_0} \left[ \hat{y} \int_0^h \left( \frac{dz}{a(a^2+z^2)^{1/2}} - \frac{z^2 dz}{a(a^2+z^2)^{3/2}} \right) + \hat{z} \frac{1}{(a^2+z^2)^{1/2}} \Big|_0^h \right]$$

$$= \frac{\rho}{4\pi\epsilon_0} \left[ \hat{y} \int_0^h d \left( \frac{z}{a(a^2+z^2)^{1/2}} \right) + \hat{z} \frac{1}{(a^2+z^2)^{1/2}} \Big|_0^h \right]$$

$$= \frac{\rho}{4\pi\epsilon_0} \left[ \hat{y} \frac{z}{a(a^2+z^2)^{1/2}} + \hat{z} \frac{1}{(a^2+z^2)^{1/2}} \right] \Big|_0^h$$

$$\vec{E} = \frac{\rho}{4\pi\epsilon_0} \left[ \hat{y} \frac{h}{a(a^2+h^2)^{1/2}} + \hat{z} \left( \frac{1}{(a^2+h^2)^{1/2}} - \frac{1}{a} \right) \right]$$

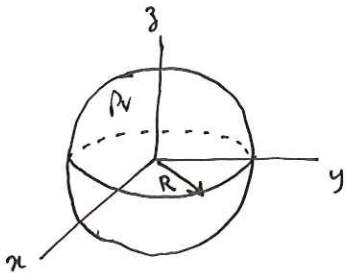
### Example 2

Given  $\vec{D} = \hat{x} 2(x+y) + \hat{y} (3x-2y)$ , find  $\rho_v$ .

Solution:  $\vec{\nabla} \cdot \vec{D} = \rho_v \rightarrow \rho_v = \frac{\partial}{\partial x} [2(x+y)] + \frac{\partial}{\partial y} (3x-2y) = 2 - 2 = 0$

### Example 3

Find the electric field around a sphere of radius  $R$  with charge density  $\rho_v$ .

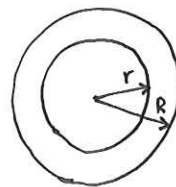


Solution: if  $r < R$ :

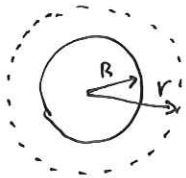
$$\int \vec{D} \cdot d\vec{s} = Q = \int \rho_v dv$$

$$D(4\pi r^2) = \rho_v \left( \frac{4}{3} \pi r^3 \right)$$

$$\vec{D} = \rho_v \frac{r}{3} \hat{R}$$



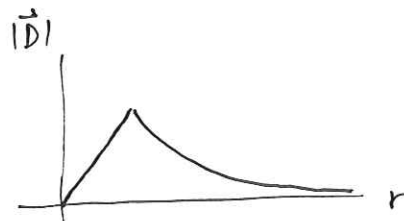
if  $r > R$ :



$$\int \vec{D} \cdot d\vec{s} = Q = \int \rho_v dv$$

$$D(4\pi r^2) = \rho_v \left( \frac{4}{3} \pi R^3 \right)$$

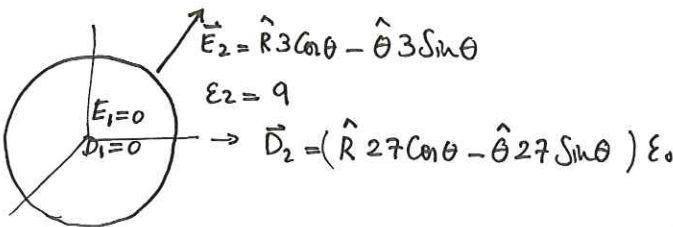
$$\vec{D} = \rho_v \frac{R^3}{3r^2} \hat{R}$$



Example 4

of radius 2cm

A spherical metal ball is inside oil with  $\epsilon_r = 9$ . If there exists an electric field of  $\vec{E} = \hat{R} 3 \cos \theta - \hat{\theta} 3 \sin \theta$  ( $\frac{V}{m}$ ) in the oil, find the charge density on the ball.



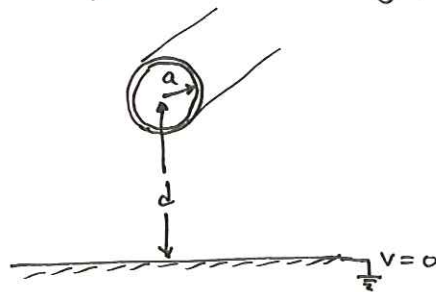
$\vec{E}_2 = \hat{R} 3 \cos \theta - \hat{\theta} 3 \sin \theta$   
 $\epsilon_2 = 9$   
 $\vec{D}_2 = (\hat{R} 27 \cos \theta - \hat{\theta} 27 \sin \theta) \epsilon_0$   
 $E_1 = 0$   
 $D_1 = 0$

$\nabla \cdot \vec{D} = \rho \rightarrow D_{1n} - D_{2n} = \rho_s \rightarrow \rho_s = -D_{2n} = -\vec{D}_2 \cdot \hat{R} = -(\hat{R} \cdot \hat{R} 27 \cos \theta - \hat{\theta} \cdot \hat{R} 27 \sin \theta) \epsilon_0$

$\rho_s = -27 \cos \theta \epsilon_0 \quad \left(\frac{C}{m^2}\right)$

Example 5

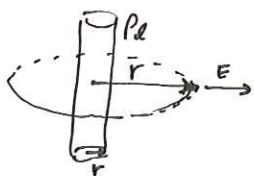
Find the capacitance per unit length of an infinitely long conducting cylinder of radius  $a$  situated at a distance  $d$  from a conducting plate as shown here.



Solution: Using the image method, we have:

$$C = \frac{Q}{V} = \frac{Q}{-\int \vec{E} \cdot d\vec{l}}$$

For a cylinder we have:



$$\int D \cdot d\vec{s} = Q = \int \rho_l \, dl$$

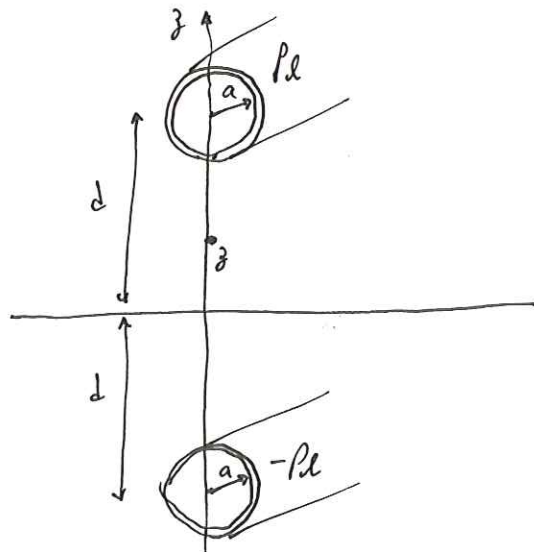
$$D \cdot 2\pi r \, dl = \rho_l \, dl \rightarrow \vec{D} = \frac{\rho_l}{2\pi r} \hat{r}$$

So for the top conductor,  $\vec{E}$  at  $z$  is:

$$\vec{D}_1 = -\hat{z} \frac{\rho_l}{2\pi(d-z)} \quad \text{and for the bottom conductor:} \quad \vec{D}_2 = -\hat{z} \frac{\rho_l}{2\pi(d+z)}$$

So the total  $\vec{D}$  at  $z$

is:



$$\vec{D}(z) = \vec{D}_1 + \vec{D}_2 = -\hat{z} \frac{\rho_l}{2\pi} \left( \frac{1}{d-z} + \frac{1}{d+z} \right) \rightarrow \vec{E} = -\hat{z} \frac{\rho_l}{2\pi\epsilon_0} \left( \frac{1}{d-z} + \frac{1}{d+z} \right)$$

So  $V$  between the two conductors is:  $V = - \int_{z=d-a}^{z=-(d-a)} E \cdot dl$

$$\rightarrow V = \frac{\rho_l}{2\pi\epsilon_0} \left( \int_{d-a}^{d-z} \frac{dz}{d-z} + \int_{d+z}^{d-a} \frac{dz}{d+z} \right) = \frac{\rho_l}{2\pi\epsilon_0} \left( -\ln(d-z) + \ln(d+z) \right) \Big|_{d-a}^{-d-a}$$

$$= \frac{\rho_l}{2\pi\epsilon_0} \ln \left( \frac{d+z}{d-z} \right) \Big|_{d-a}^{-d-a} = \frac{\rho_l}{2\pi\epsilon_0} \left( \ln \frac{2d-a}{a} - \ln \frac{a}{2d-a} \right)$$

$$V = \frac{\rho_l}{\pi\epsilon_0} \ln \frac{2d-a}{a}$$

Also:  $Q = \rho_l L \rightarrow C = \frac{Q}{V} = \frac{\rho_l L}{\frac{\rho_l}{\pi\epsilon_0} \ln \frac{2d-a}{a}} \rightarrow \frac{C}{L} = \frac{\pi\epsilon_0}{\ln \frac{2d-a}{a}} \quad \left( \frac{F}{m} \right)$

Example 6 A 20-turn rectangular coil of  $30\text{cm} \times 10\text{cm}$  is placed as shown in the picture in  $yz$  plane. If the coil which carries a current  $I=10\text{A}$  is in a magnetic flux density

$\vec{B} = 2 \times 10^{-2} (\hat{x} + \hat{y}) \text{ T}$  determine the torque on the coil.

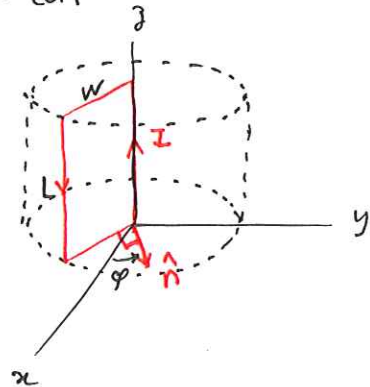
Solution: we know  $\vec{T} = \vec{m} \times \vec{B}$

$$\vec{m} = \hat{n} N I A = \hat{n} (20)(10)(0.3 \times 0.1) = \hat{n} 6$$

$$\rightarrow \vec{T} = 6 \hat{n} \times (\hat{x} + \hat{y}) (2 \times 10^{-2})$$

$$= 0.12 (\hat{n} \times \hat{x} + \hat{n} \times \hat{y}) = 0.12 (-\hat{z} \sin\varphi + \hat{z} \cos\varphi)$$

$$\vec{T} = 0.12 (2 \cos\varphi - \sin\varphi) \hat{z}$$

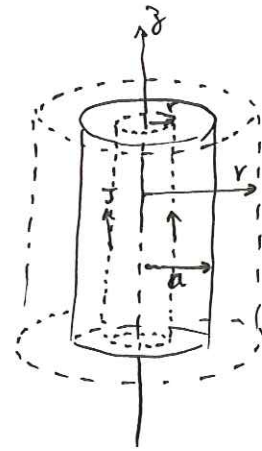


\* At what  $\varphi$ , the torque is maximum?

$$\frac{dT}{d\varphi} = 0 \rightarrow -2 \sin\varphi - \cos\varphi = 0 \rightarrow \tan\varphi = -\frac{1}{2} \rightarrow \varphi = -26.57^\circ \text{ or } 153.43^\circ$$

Example 7 Find the magnetic field around a cylindrical conductor carrying current

density  $\vec{J} = \hat{j} J_0 e^{-r}$  as shown for  $r < a$ .



$$\oint_C \vec{H} \cdot d\vec{l} = I = \int_S \vec{J} \cdot d\vec{s}$$

$$\begin{aligned} H(2\pi r) &= \int_0^r J_0 e^{-r} 2\pi r dr = 2\pi J_0 \int_0^r r e^{-r} dr \\ &= 2\pi J_0 \left( -r e^{-r} + \int e^{-r} dr \right) \Big|_0^r \\ &= 2\pi J_0 \left( -r e^{-r} - e^{-r} \right) \Big|_0^r = 2\pi J_0 \left( -r e^{-r} - e^{-r} + 1 \right) \end{aligned}$$

$$\Rightarrow H = J_0 \left( \frac{1}{r} - e^{-r} - \frac{1}{r} e^{-r} \right) \quad r \leq a$$

for  $r > a$ :

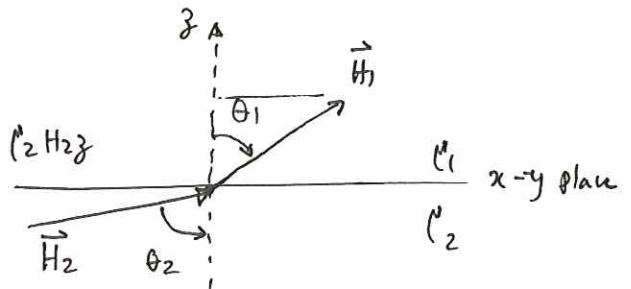
$$H(2\pi r) = \int_0^a J_0 e^{-r} 2\pi r dr = 2\pi J_0 \left( 1 - a e^{-a} - e^{-a} \right)$$

$$\rightarrow H = \frac{J_0}{r} \left( 1 - a e^{-a} - e^{-a} \right) \quad r \geq a$$

Example 8 In the picture  $\vec{H}_1 = \hat{x} H_{1x} + \hat{y} H_{1y} + \hat{z} H_{1z}$  and there is no surface current at the interface. Find  $\vec{H}_2$ ,  $\theta_1$ ,  $\theta_2$ .

$$\nabla \cdot \vec{B} = 0 \rightarrow B_{1n} = B_{2n} \Rightarrow \mu_1 H_{1n} = \mu_2 H_{2n} \rightarrow \mu_1 H_{1z} = \mu_2 H_{2z}$$

$$\nabla \times \vec{H} = 0 \rightarrow H_{1t} = H_{2t} \rightarrow \begin{cases} H_{1x} = H_{2x} \\ H_{1y} = H_{2y} \end{cases}$$



$$\rightarrow \vec{H}_2 = \hat{x} H_{1x} + \hat{y} H_{1y} + \hat{z} \frac{\mu_1}{\mu_2} H_{1z}$$

$$\tan \theta_1 = \frac{H_{1t}}{H_{1n}} = \frac{\sqrt{H_{1x}^2 + H_{1y}^2}}{H_{1z}} \quad \tan \theta_2 = \frac{\sqrt{H_{1x}^2 + H_{1y}^2}}{H_{1z}} \quad \frac{\mu_2}{\mu_1} = \frac{\mu_2}{\mu_1} \tan \theta_1$$

$$\rightarrow \frac{\tan \theta_1}{\mu_1} = \frac{\tan \theta_2}{\mu_2}$$